NACA RM A54C2

To



RESEARCH MEMORANDUM

EFFECTS OF AIRFOIL PROFILE ON THE TWO-DIMENSIONAL

FLUTTER DERIVATIVES FOR WINGS OSCILLATING

IN PITCH AT HIGH SUBSONIC SPEEDS

By John A. Wyss and James C. Monfort

Ames Aeronautical Laboratory
Moffett Field, Calif.

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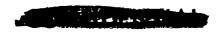
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NATIONAL ADVISORY COMMITTEE FOR AERONAUTICS

WASHINGTON

May 24, 1954





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RESEARCH MEMORANDUM

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SUMMARY

Aerodynamic lift and moment flutter derivatives were determined at high subsonic speeds for a series of two-dimensional airfoils varying in thickness and thickness distribution. The wings were sinusoidally oscillated about the quarter-chord axis at Mach numbers from about 0.5 to 0.9. The corresponding reduced frequency ranges varied from 0.045 to 0.45 at M = 0.5 and from 0.025 to 0.25 at M = 0.9. An evaluation of the results indicated that wing profile and angle of attack have major effects on the flutter derivatives at speeds exceeding the Mach number for steady-state lift divergence. In general, at supercritical Mach numbers the trends of the magnitudes of the oscillatory lift coefficients were qualitatively indicated by the trends of the nonoscillatory coefficients, with phase angles, except for the 12-percent-thick airfoil, having only moderate deviation from subsonic theory. The variations in the magnitude of the moment derivative and in its phase angle, resulted in a trend toward instability at supercritical Mach numbers. In particular, for airfoils of equal thickness the effect of an extreme forward location of maximum thickness was destabilizing in that negative -aerodynamic damping existed, implying the possibility of a single degree of freedom type of flutter. Decreasing airfoil thickness delayed the large deviation from subsonic theory to higher Mach numbers.

INTRODUCTION

This report is concerned with the evaluation of the effects of airfoil profile on the lift and moment flutter derivatives as measured, by means of pressure cells, on harmonically vibrating two-dimensional wings at high subsonic speeds. It is well-known that theory does not account properly for such factors as flow separation and shock formation, hence, the aircraft designer must of necessity look to experimental values



whenever such mixed-flow conditions are encountered. Numerous previous investigations at lower speeds, such as those by Clevenson and Widmayer (ref. 1) and by Halfman (ref. 2), may be cited. With the use of a different measuring technique, the present work extends these previous investigations to higher Mach numbers so that emphasis may be placed upon supercritical speeds for which information is meager or nonexistent.

Since wing profile may be expected to have a significant effect on mixed-flow conditions, several models were used to determine the effects of wing thickness and thickness distribution on the flutter derivatives. NACA 65A series symmetrical airfoils, 12, 8, and 4 percent thick, were used along with two other 8-percent-thick airfoils with their maximum thickness at about 16 and 63 percent of the wing chord. The models were oscillated about the quarter-chord axis at Mach numbers from 0.5 to 0.9 with reduced frequency ranges from 0.045 to 0.45 and from 0.025 to 0.25, respectively. Reynolds numbers, based on the airfoil chord, varied from 5 to 8 million.

SYMBOLS

a	velocity of sound in undisturbed air, ft/sec									
ъ ;	wing semichord, ft									
c ₁ .	namic section lift coefficient									
c _m	dynamic section moment coefficient about quarter point of chord									
f	frequency of oscillation, cps									
k	reduced frequency, $\frac{\omega b}{V}$									
M .	Mach number, $\frac{V}{a}$									
M_{CL}	oscillatory aerodynamic section moment on wing about axis of rotation, positive with leading edge up									
P_{α}	oscillatory aerodynamic section lift on wing, positive upwards									
q	free-stream dynamic pressure, lb/sq ft									
V	free-stream velocity, ft/sec									



α	oscillatory angular displacement (pitch) about axis of rota- tion, positive with leading edge up, radians
$\alpha_{\underline{m}}$	mean angle of attack about which oscillation takes place, deg
в	phase angle between oscillatory moment and position α , positive for moment leading α , deg
φ	phase angle between oscillatory lift and position α , positive for lift leading α , deg
(à	circular frequency, 2xf, radians/sec
$\frac{de_l}{d\alpha}$	magnitude of dynamic lift-curve slope, $\frac{P_{\alpha}e^{-i\phi}}{2bq\alpha}$, per radian
$\frac{\mathrm{d}\mathbf{c}_{\mathrm{m}}}{\mathrm{d}\alpha}$	magnitude of dynamic moment-curve slope, $\frac{M_{\alpha}e^{-i\theta}}{4b^2q\alpha}$, per radian
$\frac{dc_m}{d\alpha}$ sin θ	aerodynamic damping component in phase with angular velocity

APPARATUS AND METHOD

Models and Instrumentation

The 12- and 8-percent-thick airfoils, NACA 65A012, 65A008, 2-008, and 877A008 profiles, were of wood-rib and wood-stressed-skin construction built around steel spars at the quarter chord, which was the axis of rotation. Several wood spars at other chordwise locations were used to minimize spanwise twisting since the models were driven from one side. The 4-percent-thick model, of NACA 65A004 profile, was machined from solid aluminum with a parting line in the chord plane. The upper and lower halves of this model were bolted and doweled together. Each model had a chord of 24 inches and a span of 18-1/4 inches. The gaps between the ends of the models and tunnel walls were sealed with sliding spring-loaded felt pads or brass strips which moved with the models.

¹An NACA 847AllO airfoil was modified to a symmetrical section by using the lower surface coordinates for both upper and lower surfaces and then reducing the thickness ratio to 8 percent.



In figure 1, the model profiles are illustrated to show the variation of thickness and thickness distribution. The reference model, NACA 65A008, is marked to indicate the locations of the pressure cells. Model instrumentation consisted of 15 flush-type pressure cells (see refs. 3 and 4) and 15 pressure orifices along the midspan of each surface of each model. The pressure orifices adjacent to each pressure cell were used for two purposes: (1) as a means to determine the time-average chordwise pressure distribution with the use of a multiple mercury manometer, and (2) to provide an internal reference pressure for the pressure cells. The tubes from each cell and from the adjacent pressure orifice were interconnected at the manometer. In order that the internal reference pressure of the pressure cells would be essentially steady, about 50 feet of 1/16-inch tubing was used from the orifice to the manometer and back to the pressure cell.

Two 14-channel oscillographs were used to record the instantaneous electrical difference of the output of each pair of cells (proportional to the pressure difference between the upper and lower surface at each chord station) and to record the summation of all cells (proportional to the variation of the lift force). The output of an NACA slide-wire position transducer, proportional to the model angle of attack, was simultaneously recorded.

Tunnel, Model Drive System, and Tests

The models were oscillated in the two-dimensional test section in the Ames 16-foot high-speed wind tunnel (ref. 5). The two-dimensional channel was about 20 feet long and 16 feet high. A view of a model in place and a diagrammatic sketch of the drive system are presented in figure 2. The drive rods and sector arm attached to the model were contained within one of the channel walls.

Records were obtained with Mach number and mean angle of attack constant for frequencies from 4 to 40 cycles per second at intervals of 4 cycles per second and for an amplitude of ±1°. Data are presented for mean angles of attack of 0° and 2° and for Mach numbers from 0.5 to about 0.9. Sample oscillograph records which illustrate the necessity for harmonic analysis at the higher Mach numbers are given in figure 3. The lift was evaluated by a 12-point harmonic analysis of three consecutive cycles of the sum trace. The pitching moment was evaluated by a 12-point harmonic analysis of the individual cell traces for one cycle.

Since the investigation was conducted in a closed-throat tunnel, the effects of wind-tunnel resonance must be accounted for either by avoiding conditions in which tunnel-wall effects are significant or by correcting the results for the effects of the tunnel walls (refs. 6 and 7). Calculations made at the Langley and Ames Laboratories employing

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the single-doublet-line, single-control-point solution described in reference 7 yielded the following results for a tunnel height of 16 feet, wing chord of 2 feet, and Mach number of 0.7: At frequencies of 10, 20, and 40.66 cycles per second, the magnitudes of the coefficients were increased by 3.8, 5.0, and 4.7 percent, respectively, due to the presence of the tunnel walls. These results indicate that, for the conditions of the calculations, the effect of the tunnel walls was small. However, for mixed-flow conditions, the application of such corrections based on potential flow would be questionable; hence, to minimize tunnel-wall effects, all data obtained at frequencies within 10 percent of the tunnel resonant frequency (refs. 6 and 7) have been omitted. Although the use of such a procedure does not mean tunnel-wall effects have been completely eliminated over the entire frequency range, it is felt that tunnel-wall effects are not a predominant factor in the trends of the data.

For a discussion of other factors influencing the precision of the data, the reader is referred to references 3 and 4.

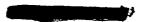
RESULTS AND DISCUSSION

A tabulation of the measured derivatives is contained in tables I, II, III, IV, and V for the NACA 65A012, 65A008, 65A004, 2-008, and 877A008 airfoils, respectively. The results concerning lift derivatives are first discussed and are presented in figures 4 to 10, followed by a discussion and the presentation of the moment derivatives in figures 11 to 15.

Lift

Experimental values for the reference model for three representative Mach numbers are presented in figure 4 as a function of reduced frequency. In this figure, as in subsequent figures, the absolute magnitude of the flutter derivative is expressed in terms of the slope of the lift curve per radian and the corresponding phase-angle relationship between the lift vector and model angle of attack in degrees. Theoretical values at Mach numbers of 0.5, 0.6, and 0.7 may be obtained from the work of Dietze (refs. 8 and 9), and at Mach numbers of 0.8 and 1.0 from Minhinnick (ref. 10) and Nelson and Berman (ref. 11), respectively.

In this figure it may be noted that at 0.49 and 0.79 Mach numbers the flutter derivatives tend to increase with increasing reduced frequency; furthermore, there seems to be a large variation in the phase angle at low values of reduced frequency at 0.79 Mach number. However,



Mach number appears to have had a greater effect on the data than did frequency at 0.91 Mach number.

Typical results as a function of Mach number are presented in figure 5 for the reference model, the NACA 65A008 airfoil. The lines showing the theoretical values are identified at one end by the frequency in cycles per second to which they pertain. Since theoretical values have been computed in the cited references only at certain Mach numbers which have already been indicated, an interpolation was necessary to obtain values at intermediate Mach numbers. Although such an interpolation inherently involves some error, a consistent set of values was nevertheless established and was used for the purpose of determining the effects of varying airfoil shape.

To distinguish between the various frequencies, the experimental and theoretical values are each faired with the same type of line. For example, the experimental and theoretical values for a frequency of 8 cycles per second are each shown with a solid line. Examination of the experimental data for a frequency of 8 cycles per second indicates that the trends of both experiment and theory were the same at low Mach numbers. As Mach number increased, a large decrease in the magnitude of the experimental derivative occurred, accompanied by a variation of phase angle such that the trend toward increasing lag was reversed. Although the agreement with theory was not precise at the lower Mach numbers, it may be seen that the general trends for all frequencies were nearly the same.

The data from figure 5 are presented in a different form in figure 6; the experimental magnitude has been divided by the theoretical magnitude, and the theoretical phase angle has been subtracted from the experimental phase angle. These quantities are also shown as a function of Mach number. If the experimental and theoretical values exactly agreed, the ratio of the magnitudes of the derivatives would be 1, while the difference in phase angle would be 0. The faired lines represent the average deviation from theory for the entire frequency range up to 40 cycles per second.

It is of interest to note that the individual points do not indicate an entirely random scatter about the mean line for the various frequencies. For example, examination of the points for 40 cycles per second in the top portion of the figure shows that these points are usually the uppermost value at each Mach number. Hence, this figure not only provides some indication of the range of the experimental values, but illustrates the fact that, although the values depend on frequency, the general variations with Mach number are represented by the faired average curves.

The use of the average deviation from theory appears to be justified since it is representative of each model. For example, in figure 6 it may be noted that all the experimental points lie within a comparatively



narrow band along the faired curves with the exception of the higher frequencies in the upper portion of the figure. In fact, a band of width ±0.15 in the upper portion of the figure and a band of width ±10° in the lower portion of the figure would contain about 80 percent of all the experimental points. These results are typical of all the models. It might be noted that the averaging process used has the effect of removing frequency as a parameter. It should be noted that each model was oscillated at the same amplitude and through the same range of frequencies, hence the average deviation from theory indicates the over-all effects of airfoil shape and the general trends of the data.

Effect of thickness distribution.— The effects of the variation of thickness distribution as indicated by the curves showing the average deviation from theory over the frequency range tested are summarized in figure 7 for mean angles of attack of 0° and 2°. It would appear from this figure that the main effect of the chordwise location of maximum thickness was on the magnitudes of the derivatives rather than on phase angles, although no systematic trend is apparent.

Effect of wing thickness. The results showing the effects of wing thickness are presented in figure 8. At an angle of attack of 0° , wing thickness appears to have had a much more pronounced effect than wingthickness distribution (fig. 8(a) as compared to fig. 7(a)). As might be expected, the primary effect of reducing wing thickness was to delay any large deviation from theory to a higher Mach number.

At an angle of attack of 2° (fig. 8(b)), large differences over the entire range of Mach numbers occurred between the models in the magnitudes of the derivatives.

Comparison with steady-state results. In order to examine whether any relation existed between unsteady and steady-state results, a comparison with steady-state results obtained from the time-average chord-wise pressure distributions for mean angles of attack of 0° and 2° is made in figures 9 and 10. In these figures, the steady-state data have been normalized with the Prandtl-Glauert value of the theoretical lift-curve slope. It may be recalled that the Prandtl-Glauert curve is also obtained as an end condition as the frequency of oscillation approaches zero.

Examination of these figures indicates that although there appears to be some parallelism or similarity between the steady and unsteady curves, the comparison between the steady and unsteady values is at best only qualitative. For example, in neither figure 9 nor figure 10 do the unsteady and steady-state curves coincide throughout the entire range of Mach numbers. It should also be noted that, with the exception of the NACA 65A012 airfoil at a mean angle of attack of 2° (fig. 10(b)), the unsteady values approached theory more closely than did the steady-state



values, particularly at the lower Mach numbers, that is, from M=0.5 to 0.7. Although the effect of the higher frequencies in increasing the level of the curves for the unsteady case may in part account for the differences between the curves, this effect is small. However, the one characteristic that is common to both the unsteady and steady curves in almost every case is a trend toward a reduction in magnitude at the highest Mach numbers. The Mach number at which this trend initiates cannot be precisely delimited, nevertheless, for the three NACA 65A-series airfoils at a mean angle of attack of 0° (fig. 10(a)), the unsteady lift trend appears to be associated with the steady-state flow changes which occur above the Mach number for lift divergence.

It would therefore appear that as a first approximation the Mach number for lift divergence may be taken as a criterion for the onset of significant changes in the trends of the unsteady values, and that this trend toward a decrease in the magnitude of the unsteady values is related to the trend of the steady-state data. It should be pointed out that this conclusion is not as evident for the NACA 2-008 and 877A008 airfoils (fig. 9) and for the NACA 65A004 airfoil at a mean angle of attack of 2° (fig. 10(b)), since these figures indicate that the correlation between the Mach number for lift divergence and the initiation of a downward trend of the unsteady values is not precise and they may differ by as much as 0.1. However, it is felt that there is sufficient evidence presented in figures 9 and 10 to indicate that steady-state values may prove useful as a qualitative indication of the trends of the unsteady-state coefficients at supercritical Mach numbers.

For the steady-state condition the phase angle is, of course, zero; therefore no corollary for the phase angle with relation to the oscillatory condition is possible. However, except for the 12-percent-thick wing, the phase angle shows only a moderate deviation from theory throughout the speed range of the present investigation.

Moment

The moment derivatives for the reference model as a function of reduced frequency for several Mach numbers are presented in figure 11 and as a function of Mach number in figure 12. A comparison of these figures indicates that even though there may have been a greater effect due to frequency on the moment derivatives than had been the case for the lift derivatives, from figure 12 it appears that the effects of Mach number are similar for all frequencies. Hence, the effects of airfoil profile are again compared on the basis of the faired average curves in figure 12 which represent the average deviation from theory over the entire frequency range.



In contrast to the lift results previously presented in figure 6, the magnitudes of the moment derivatives greatly exceeded the theoretical values, along with a much larger variation of phase angle as compared with theory. These results may be attributed to the fact that the comparison is between very small quantities in regard to the magnitude of the derivatives, since the moment is taken about the quarter-chord axis, and to small movements of the center of pressure which would be reflected in large changes of phase angle. The general trends of the results, nevertheless, are represented by the faired average curves.

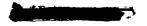
Effect of thickness distribution. The effects of the variation of the chordwise location of maximum thickness are shown in figure 13. An apparent characteristic of the NACA 2-008 airfoil, with a forward location of maximum thickness, is a large shift toward a lagging phase angle as Mach number increased above 0.8, such that the phase angle lagged theory by 80° and 90° at angles of attack of 0° and 2°, respectively. The effects of such large shifts in phase angle are discussed in relation to subsequent figures.

Effect of wing thickness. The effects of wing thickness on the moment derivatives are shown in figure 14. As might be expected, the primary effect of decreasing wing thickness was again to delay any large variations to a higher Mach number.

Instability. Since there was such a large variation at the higher Mach numbers from the subsonic theoretical values, it is of basic importance to examine the damping-moment derivatives directly to determine whether instability, or the existence of negative aerodynamic damping (implying the possibility of a single degree of freedom type of flutter), which is not predicted by the theory, existed at these speeds. The average damping-moment derivatives for the entire frequency range are therefore presented in figure 15. Also included in this figure are dashed lines indicating average values derived from theory for the corresponding frequency range.

The effect of wing-thickness distribution on aerodynamic damping is shown in figure 15(a) for each mean angle of attack. It may be noted that there was a trend toward instability for each model, with the NACA 2-008 airfoil becoming abruptly unstable at about 0.85 Mach number at 0° and 2° angles of attack. It would appear that stability about the quarter-chord axis increased as maximum thickness was moved toward the trailing edge.

The effect of wing thickness on the aerodynamic damping moment is shown in figure 15(b) for each angle of attack. Although the trend toward instability does not appear at 0° angle of attack for the NACA 65AOO4 profile, the susceptibility of the thinner wing to negative aerodynamic damping is clearly indicated at the 2° mean angle of attack.





CONCLUSIONS

Within the limitations of speed range and angle-of-attack variation of the investigation, the following general conclusions may be drawn:

- 1. Section profile has a major effect on the flutter derivatives at speeds exceeding the Mach number for steady-state lift divergence.
- 2. It appears that the variation in angle of attack has an effect as important as the effect of the variation in profile.
- 3. In general, at supercritical Mach numbers, a qualitative evaluation of the results indicated that the trends of the magnitudes of the oscillatory lift coefficients were indicated by the trends of the non-oscillatory lift coefficients, with phase angles, except for the 12-percent-thick model, showing only a moderate deviation from theory.
- 4. The variations in the magnitude of the moment derivative and in its phase angle, resulted in a trend toward instability at supercritical Mach numbers. In particular, for airfoils of equal thickness the effect of an extreme forward location of maximum thickness was destabilizing in that negative aerodynamic damping existed, implying the possibility of a single degree of freedom type of flutter.

Ames Aeronautical Laboratory
National Advisory Committee for Aeronautics
Moffett Field, Calif., Mar. 24, 1954

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TABLE I.- MEASURED FLUTTER DERIVATIVES FOR THE NACA 65A012 AIRFOIL

K 0.491	k	6	dc1		α ₁₀ = 0 ⁰								α ₁₁₁ = 2 ^O						
0.491			åα	ф	dc _m	6	м	k.	3	dc 1	ф	dc _m	6						
	0.103 .184 .282	57:0 101.8 155:9	6.394 5.466 5.099	351.8 358.8 355.5			0.491	0.058 .094 .136 .187	31.7 51.1 74.1 102.5	6.520 5.578 5.574 4.989	354.6 354.1 0.0 5.3								
.590	.077 .152 .229	51.6 101.3 153.2	7.083 6.056 5.319	351.7 351.9 351.2				.238 .287 .328 .469	130.1 157.1 179.5 256.5	4.987 5.341 5.058 4.823	4.5 0.0 12.4 29.4								
.633	.074 .111 .144 .183 .218 .252 .320	52.6 79.2 103.0 130.9 155.9 180.0 228.5 256.5	5.745 5.068 5.299 4.661 4.449 4.036 3.913 4.259	355.0 355.5 355.5 357.3 358.0 349.6 15.0	0.531 .590 .585 	342.1 317.6 305.0 311.6	.590	.048 .076 .120 .152 .198 .233 .347	31.5 50.5 79.4 100.8 131.2 154.0 229.3 254.4	6.523 6.262 5.925 5.965 5.488 5.213 4.744 5.426	352.4 345.5 347.7 354.2 352.5 347.7 9.2 16.1	0.559 .771 .739	3 ⁴ 1.7 317.8 297.5						
.682	.064 .097 .130 .163 .197 .264 .293	49.8 76.0 101.6 127.3 153.7 206.2 229.3 254.4	7.918 7.332 6.855 5.533 5.765 4.362 5.118 4.932	344.4 339.7 348.2 337.3 346.4 2.5 0.4	.595 .658 	325.4 310.6 279.7 291.5 278.8	.682	.044 .066 .101 .131 .163 .196 .294	34.6 52.0 80.2 103.7 128.9 154.8 232.7 253.7	6.216 5.833 5.506 5.224 5.055 4.528 4.290 4.329	354.5 349.2 347.3 0.4 348.8 342.7 15.3 2.0								
.731	.062 .098 .121 .156 .247 .280 .308	51.2 81.2 100.6 129.4 205.6 232.7 256.1	8.080 8.454 7.092 6.092 5.187 5.299 5.018	348.1 339.5 339.5 328.9 356.2 355.2	.634 .675 .647	326.9 304.6 283.0 279.5	.731	.041 .060 .093 .122 .153 .240	34.6 51.2 79.2 104.5 130.4 204.9 231.3	6.788 6.050 5.566 5.437 5.280 4.182 4.375	351.2 348.9 349.8 351.9 346.3 359.8 0.6	.698 .642 .721	340.5 333.2 315.6 300.4						
· .790	.057 .086 .114 .142 .199 .226	52.2 77.8 103.9 129.3 180.9 205.3 232.7	8.576 8.362 7.476 6.137 4.771 4.588 5.285	343.5 337.9 336.4 327.1 351.7 348.9 356.6	.242 .045 	305.4 276.3 263.1	.790	.034 .056 .086 .115	30.9 50.8 77.6 103.5 125.1	6.377 5.981 7.353 6.628 5.099	353.8 347.9 343.3 341.9 333.5	.597 .606 .688	292.6 340.3 316.4 277.1						
.837	.052 .077 .104 .182 .207	50.8 74.7 101.5 177.3 200.9	4.894 4.590 4.780 3.515 3.597	354.0 342.7 351.6 2.6 12.2	.612 .828 .857	301.6 269.9 256.9	.837	.198 .225 .254 .279	178.5 202.7 228.8 251.6	3.861 4.047 4.196 4.895 4.318	353.3 348.9 358.7 0.9	.557	287.3 281.6 340.8						
.885	.235 .262 .030 .049 .080	228.4 255.4 30.7 50.3 82.1 99.7 153.2	965 .641 1.725 1.884 2.681	16.5 359.9 47.0 92.7 59.9 47.4 41.2	1.736 2.719 3.117 2.436 1.939	222.0 348.1 348.0 311.4 314.4		.054 .080 .103 .181 .208 .244 .261	51.8 77.7 100.2 175.5 201.7 236.5 252.3	4.585 4.775 4.570 3.654 4.068 4.601 5.379	357.6 353.3 356.9 348.1 351.4 13.9	. 675	306.3 281.4 259.2 242.9						
	.176 .201 .223 .246	181.3 207.3 230.4 253.3	2.015 1.454 2.733 2.681	29.0 33.2 32.1 2.4	1.223	304.0	.885	.030 .049 .080 .097 .149 .177 .202 .225	30.7 50.3 82.1 99.7 153.2 181.3 207.3 230.4 253.3	3.497 2.751 3.032 2.403 2.647 3.564 2.084 3.944	347.1 345.4 359.2 342.9 353.5 351.6 336.5 345.6 356.4	1.256 .879 .712 1.389 1.043	356.8 350.2 340.7 0.2 333.5						





TABLE II.- MEASURED FLUTTER DERIVATIVES FOR THE NACA 65A008 AIRFOIL

			α <u>π</u> = 0°				α ₂₄ = 5 ₀							
и	Ŀ	ω	de ₂	φ	dc _a	8	ж	k	a)	đe z	φ.	dc _u	8	
0.491	0.089 .142 .184	48.9 78.2 101.3	6.186 5.638 5.507	354-7 348-3 353-3			0.492	0.056 .089 .140	30.7 48.9 76.7	5.356 5.074 4.613	352.1 354.4 351.5		===	
	.234 .280 .322 .457	128.5 153.9 177.0 251.3	5.319 6.250 5.518 5.883	357.0 1.9 352.5 19.2				.231 .281 .323	101.8 126.7 154.4 177.5 254.4	4.571 4.205 4.510 4.443 4.828	357.1 358.0 4.6 356.0 24.2			
.590	.074 .116 .152 .191 .234	19.2 17.0 101.2 127.2 155.9 179.5	6.657 5.841 5.756 5.854 5.673 5.500 5.612	351.0 343.6 344.9 347.5 349.1 345.2	0.445 .626 .775	315.8 312.2 289.8	.590	.031 .076 .117 .153	20.3 50.6 77.8 102.0 125.7	6.114 5.803 5.263 5.237 4.883	349.6 344.2 346.4 348.0	0.581 -557 -588	341.1 331.6 307.1	
.680	.347 .383	231.0 254.4 50.1	6.787	6.7 9.5 348.9 343.0	1.213	290.0		.231 .271 .346 .378	153.2 180.0 230.1 251.3	4.975 4.646 4.122 5.246	346.9 345.5 13.9 15.7	.701	281.4	
	.102 .130 .167 .199 .295 .329	79.4 101.0 129.3 154.4 228.8 255.1	6.312 6.430 6.062 5.652 6.067 6.389	343.0 338.4 335.7 336.2 2.0 2.3			.680	.038 .066 .099 .133 .165	29.3 51.2 76.3 102.6 127.9	6.618 6.183 5.972 5.753 5.641	351.7 346.8 343.3 341.5 335.9	.778 .806 .811	341.0 331.4 309.7	
.728	.060 .092 .123	50.4 77.2 102.5	7.392 7.005 6.535	340.9 339.5 335.5			en d	.200 .296 .327	154.4 228.8 252.6	5.339 5.392 5.662	330.3 357.5 350.6	1.097	278.9	
	.154 .245 .276 .305	129.1 204.9 230.4 254.9	6.028 5.696 6.297 6.196	331.0 347.0 348.9 352.4			.728	.037 .063 .095 .127 .156	30.7 52.8 80.4 106.7	7.311 7.347 6.759 6.467 6.247	350.6 343.1 337.6 338.7 326.4	.888 .944 .972	350.5 330.8 304.7	
.786	.058 .086 .114 .143	52.3 78.7 103.9 130.2	7.999 7.381 6.851 6.132	339.5 336.1 326.9 320.9 348.3	.837 .805	323.9	763	.215 .277 .303	131.3 206.2 233.0 255.4	5.508 6.156 5.911	349.3 347.3 352.0	.979 1.233	292.6 277.5	
	.225 .252 .279	181.4 204.9 229.9 254.4	5.394 5.525 5.800 6.848	344.5 347.7 345.7	.829 1.367	284.5 271.4	.761	036 059 117 126	31.8 51.8 80.0 102.9 128.9	6.523 7.863 6.883 6.377 6.005	341.9 341.2 331.2 320.8			
.833	.050 .080 .105 .214 .238	47.7 76.1 100.7 204.9 227.6	7.488 6.865 6.343 4.705 5.424	335.9 332.5 325.5 356.2 353.5	.454 .487 .651	263.8 291.9		.296 .234 .261	182.0 206.0 229.8 255.5	5.224 5.879 6.018	351.8 350.2 346.7 354.4			
.879	.263 .026	251.3 26.7 49.2	6.365 9.006 6.862	345.7 336.4 349.4	.223 .203	272.6 198.6 281.9	.786	.034 .057 .084 .114	31.1 52.5 76.7 104.1	9.568 8.362 7.931 7.520	347.1 337.1 333.0 326.8	1.086 1.092 1.053	337.5 320.2 286.4	
	.074 .149 .174 .200 .223 .248	75.7 152.9 178.8 205.6 229.0 255.1	7.193 4.299 4.315 5.360 5.990 6.383	327.6 353.4 358.8 348.1 337.5	.093 .517	155.9 245.9 189.3		.144 .199 .225 .254 .263	131.6 181.8 205.8 232.1 258.2	6.300 4.694 5.278 5.627 6.853	318.2 349.1 340.5 343.6 343.5	.973 1.526	275.8	
-917	029 048 071	31.5 51.3 _79.0	3.742 3.360 2.973 3.437	349.8 344.3 341.9	1.157	327.8 319.3	.833	.030 .055 .061 .107	29.4 53.5 79.4 104.4 155.6	8.034 8.015 7.464 6.595 4.320	347.4 338.4 330.0 325.5 344.1	.487 .532 .479	341.9 302.9 254.9 284.9	
	.145 .167 .188 .214	156.3 180.3 202.4 230.5	3.654 3.634 3.515 4.228	346.7 348.4 344.9 346.1 349.5	.922	300.6 271.6		.186 .213 .236 .262	181.1 208.3 230.7 256.1	5.163 5.302 6.237 5.502	318.2 355.2 351.0 310.8	.678 1.013	286.3	
	.236	254.3	3-997	341.9	1.440	252.4	.879	.031 .053 .075	31.7 55.2 77.7	7.030 7.460 6.721 5.835	347.9 331.6 333.8 330.1	.380 .294	311.2 284.4	
								.099 .150 .174 .199 .225	103.2 155.9 180.7 206.5	4.641 4.870 6.217	336.5 351.0 348.6 344.1	.323 .197 1.062	319.9	
					i			.251	233.5 261.1	6.666 7.154	337-3	1.175	212.0	



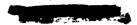


TABLE III. - MEASURED FLUTTER DERIVATIVES FOR THE NACA 65A004 AIRFOIL

	α _m = 0 ⁰						α _m = 2°						
ж	k	ω	ga	φ	do _m	θ	Ж	k	tu)	da da	φ	de _E	6
0.594	0.040 .080 .109	27.1 53.4 73.1	6.212 5.849 5.411	355.3 350.1 357.0	0.466 .447	334.7 320.4	0.491	0.046 .095 .134	25.5 52.5 74.0	6.069 5.699 5.441	357.7 2.6 5.4	= = =	
	.149 .188	99.9 125.7	5.186 5.195	358.3 8.2	.517	308.1		.186	102.5	5.176 4.848	8.2	===	= = =
	.224 .261 .342	153.3 178.0 233.5	4.883 5.242 6.989	353.5 353.0 353.8	.771	284.8		.267 .309 .448	150.3 174.0 252.3	4.367 4.702 9.199	12.4 2.0 341.4		===
.691	.382	260.9	8.995 6.846	330.7 356.1	1.396	219.8	.590	.040 .080	26.9 53.0	5.862 6.011	355.7 357.3	0.656	3 4 3.0 33 4. 3
, , , ,	.06 ¹ 4	50.0 75.5	6.669 6.039 5.887	350.8 349.9	- - -	= = =		.112 .154	74.3	5.873 5.362 5.091	359.8	.859	331.2
	.131 .162 .194	101.9 126.1 153.9	5.459 5.251	351.2 355.9 340.4	= = =			.193 .226 .258 .338	131.4 154.0 176.5	5.501 6.340	351.8 351.6 347.7	.790	287.4
	.258 .294 .327	204.9 233.5 259.6	4.981 7.405 7.473	3.0 346.9 324.3				.338 .370	231.0 252.3	7.528	2.7	1.282	256.7
.741	.032	26.6 54.1	7.366 6.996	354.9 349.7	.561 .578	338.2 327.8	.691	.033 .069 .097	25.6 53.2 74.6	7.400 6.966 6.718	355.6 351.9 356.6	.808 .839	341.0 330.2
	.089 .117	77.3	6.528 6.235 5.861	347.5 345.9 339.7	.629	301.1		.133 .164	102.4	6.316 6.397 6.526	5.6 356.4	.856 .984	332.8
	.169 .239	126.5 145.3 205.3	5.078 5.599	329.0 357.6	.607 .772	252.3 272.3		.192 .285 .321	153.0 227.1 255.8	8.342 9.530	339.5 350.1 330.7	1.564	243.7
	.269 .303	231.3 260.7	7.073 7.035	346.5 337.8	1.136	285.1	.741	.030	25.2 49.3	8.117 7.560	353.9 351.7	.975 .970	345.8 333.5
.798	.030 .056 .083	27.4 50.2 75.1	7.739 7.522 7.157	354.6 345.7 343.9	.609 .658	340.3 320.5		.087 .120 .147	72.5 100.0 122.7	7.252 6.935 6.574	353.2 358.4 359.1	1.004	325.6
	.115	103.7 127.6 183.7	6.606 5.950	343.5 335.7 343.5	.705	292.7		.175	151.6 202.7 224.7	6.128 5.990	332.8 354.9	1.034 .840	276.5 287.3
	.197 .218 .250	203.0 233.0	5.301 5.477 7.536	343.2 343.4	.864 1.217	274.3 279.4		.260 .297	256.8	8.895 8.784	339.4	1.536	244.6
.850	.278	259.1 25.5	7.924 9.637	323.1 353.9 344.4			.798	.029 .059 .083	26.0 53.5 74.5	9.165 9.012 8.581	352.9 347.4 351.7	1.140	344.5 327.2
	.053 .078 .106	50.9 74.8 102.2	8.200 7.611 7.025	344.4 340.3 337.4 347.2	.664 .736	313.6 287.1		.114 .135 .191	102.5 122.0 178.8	7.322 6.416 6.260	349.1 346.1 349.0	1.209	318.2
	.153 .179 .208	153.4 178.7 208.0	4.590 4.602 5.178	347.2 356.0 351.7	.588 .928	276.5		.218 .244 .275	204.0 228.5 258.2	7.074 9.002 10.446	343.7 339.2 335.9	2.054	280.1
	.231 .259	231.4 258.5	6.964 8.436	298.2 325.4	1.611	251.9	.852	.029	25.6	10.793 9.672	348.5 346.1	.966	334.7 321.2
.900	.025	25.4 51.1	7.315 8.913	347.3 344.6	.717 .738	331.8 310.8		.086	52.3 75.3 103.0	9.016	339.5 345.2	1.364	301.6
	.073 .126 .147	74.1 128.2 155.4	8.483 5.365 4.721	337.4 329.0 341.9	.452	270.6		.139 .176 .204	121.8 177.5 205.1	6.564 5.269 5.928	326.2 345.7 353.7	1.430	282.7
	.170 .195 .220	180.6 206.6 232.7	4.819 8.097 8.076	352.4 0.6 337.7	1.208	281.8 305.2		.228 .253	230.2 254.7	8.365 9.206	336.3 338.2	2.450	279.5
alia	.244	258.2	7.635	323.5			.870	.026	49.7	11.945	344.7 339.1		
.942	.025 .050 .096	26.1 53.2 102.2	5.877 9.448 6.885	344.2 335.5 316.1	.438 .824 .619	328.4 290.3 298.1		.074 .101 .142	74.7 101.7 151.4	9.377 7.206 4.826	327.6 322.6 342.6	.¥05 .510 .311	246.6 229.7 341.8
	.120	127.4	4.964	334.7				.164 .190 .212	174.9 202.7 226.4	5.713 6.913 7.880	356.1 350.5 340.0	.908	288.9
							.904	.238	254.7	9.150	331.8	1.546	215.5 180.1
							.,,,,	.050	72.7 72.1	9.216	326.0 325.6	1,162	147.2
								.117	123.5	5.181	336.7		





TABLE IV. - MEASURED FLUTTER DERIVATIVES FOR THE NACA 2-008 AIRFOIL

<u> </u>			α _m = 0 ^c			-	α _m = 2°						
М	[dc-1] [dcm]						ж	k	(4)	dc 1	ф	dc _m	6
0,590	0.040 .081 .113 .155 .193 .229 .350	26.3 53.9 75.3 103.0 128.7 152.5 232.7 259.6	6.460 6.082 5.624 5.436 5.425 5.287 5.001 6.210	354.6 353.5 351.4 352.3 351.2 351.2 7.0 1.2	0.552 .591 .620 .705	347.9 327.2 309.3 288.4 263.9	0.491	0.052 .093 .137 .181 .228 .278 .327 .465	28.6 51.5 75.9 100.2 126.2 154.0 181.1 257.5	6.240 5.919 5.838 5.329 5.324 5.543 5.396 5.638	358.2 355.0 351.4 355.5 356.2 1.2 3.1 22.5		
.680	.036 .069 .098 .134 .164 .197 .266 .299	27.9 53.0 75.5 103.9 127.2 152.4 206.0 231.0 254.7	7.006 6.433 6.118 5.600 5.437 5.227 4.627 5.723 5.931	351.6 347.9 349.0 347.1 346.0 343.6 358.3 356.1 350.6	.538 .615 .613 .692 .587	329.3 329.1 306.9 284.1 275.7	•590	.040 .080 .113 .146 .193 .231 .346 .384	26.8 53.9 75.6 98.2 129.6 155.1 231.9 257.5	6.691 6.204 6.093 5.706 5.650 5.665 5.379 6.678	354.2 346.8 348.4 352.9 349.6 348.5 4.8 1.4	0.574 .639 .604 .737	335.9 326.9 297.7 291.4 274.3
.728	.031 .063 .093 .123 .151 .249 .279	25.6 52.8 77.6 102.6 126.5 208.7 234.4 258.2	8.587 6.584 6.363 5.600 5.194 4.738 5.550 5.839	350.5 347.4 341.7 344.0 344.7 356.3 355.4 353.2	.788 .720 .738 .375	349.0 325.8 303.9 248.6 253.1	.680	.033 .069 .098 .134 .167 .200 .299	25.9 53.9 76.2 104.1 130.4 155.8 232.7 258.2	7.347 6.880 6.363 6.006 6.026 5.475 5.728 6.592	355.2 350.0 347.1 349.9 347.1 345.0 358.5 355.5	.661 .676 .645 .669	335.6 330.7 305.5 277.2 274.7
.786	.029 .056 .083 .111 .140 .196 .227	26.1 51.1 75.5 101.7 127.5 178.5 206.6 231.0	8.206 7.444 7.015 6.296 5.053 3.884 4.799 5.160	348.5 341.7 336.2 335.1 335.1 351.8 342.8 348.3	.665 .578 .659 	338.1 315.3 285.2 256.9	.728	.031 .062 .090 .122 .152 .245 .279	26.0 52.5 75.7 103.0 127.9 206.7 235.6 258.2	7.829 7.517 6.957 6.932 6.624 6.264 7.365 7.482	358.2 354.5 353.5 344.6 2.6 358.2 358.9	.642 .441 .599	341.4 301.3 286.4 269.5
.833	.280 .054 .080 .108 .164 .184 .212	255.4 26.8 52.2 77.6 104.9 159.9 179.5 206.5	9.150 8.806 7.986 6.952 4.893 5.420 5.576	349.1 344.7 337.6 332.1 327.3 338.8 337.5 346.4	.399 .245 -357 .268 	246.1 210.5 218.9 233.0 222.6	.786	.029 .056 .083 .113 .140 .227 .256 .281	26.2 51.8 76.4 104.2 129.1 208.5 235.6 258.2	8.917 8.673 8.193 7.768 6.988 5.927 6.469 8.547	348.2 344.7 342.4 341.1 330.8 343.6 344.7 347.3	.086 .098 .260 .405	320.2 296.0 253.5 263.1 258.1
.879	.239 .263 .027 .051 .076 .100 .150	232.7 256.1 27.6 53.1 78.2 103.0 155.3 177.5	6.807 7.713 9.720 8.316 7.062 5.235 3.831 4.760	348.8 337.3 348.8 332.1 325.5 324.2 340.4 312.5	.525 1.250 1.353 .986 .454	214.1 156.3 138.2 139.5 135.3	.833	.053 .078 .107 .185 .214 .239 .260	27.1 76.9 104.7 181.8 207.6 234.2 257.2	9.455 9.008 8.618 7.950 5.425 5.662 7.075 8.568	353.9 345.4 338.8 327.7 342.0 349.8 349.6 341.2	.317 .316 .339 .346 .467	203.3 200.2 172.2 261.4 281.4 293.6
	.200 .226 .252	206.0 233.3	6.218 6.432 7.107	341.4 335.2 330.6	.642 1.007	165.4 157.7	.879	.027 .052 .075 .100 .148 .174 .201 .225 .248		10.660 8.749 6.975 5.550 4.421 4.481 6.225 6.522	343.1 332.4 326.0 326.9 335.5 357.0 348.7 342.3 331.1	1.437 .527 .485 .743	144.1 96.8 179.0

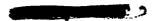
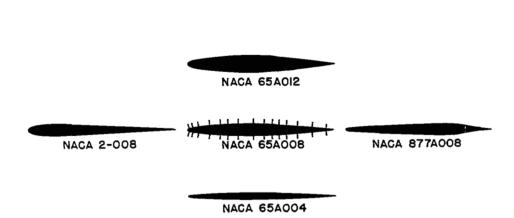




Table v.- measured flutter derivatives for the NACA 877A008 Airfoll

o ²⁸ = 0 ₀						α _m = 2°							
ж	k	u	<u>de</u> 1 de	•	da da	•	×	k	•	de l	ф	dom	•
0.195	0.050	27.9 51.5	6.087	352.7 349.5	===		0.496	0.048	27.1	6.416	353.0		
-	.092	75.7	5.638	347.1	1 = = =		1	.087	49.1 72.1	5-936 5-497	350.8 351.9	2 2 2 1	
	.183	103.0	5.159	1 3M3.2				178	100.0	5.162	349.7		
	.228	128.2	4.941	355.0		1:::		-223	125.4	5,429	347.8 344.4		
	.266	179.5	5-135	355.0 350.9 338.1				.269	151.4	4.985	352.8	===	
	.316 .450	179.5 255.4	4.486	4.8		<i>-</i>		.321 .454	253.3	5.243 4.369	18.3		
.596	.040	27.2	6.246	348.0 346.9	0.461	339.8 324.1	.596	.035	24.0	6.579 6.457 5.948	1.7	0.634	342.4
- 1	.073	19.8 74.3	5. 62	346.5	.430	324.1		.076	51.3 73.6	5.048	350.6 348.7	.621	333.5
- 1	.153	103.7	5.448	345.2	-609	310.6	}	.147	99.6	5.642	350.8	-765	320.9
- 1	.187 .223	126.9 151.7	5.370 4.728	345.5 346.0	.560	301.9		.182	123.9	5.798 5.087	352.7 347.1	.625	213.6
- 1	•336 •375	231.0	4.267	355.1				.259	174.5	5.325	340.3		
1	-375	257.5	5.053	3-4	.963	293.8	1	•337	227.6	5.325 4.688 5.319	357.6 3.6	1.192	197.6
.693	.035	27.5	6.976 6.842	348.5 343.4	= = =	:::	.693	.372		1			
i	.096	76.7	6.423	3+3.0			.093	.031 .064	25.1 51.3	6.333	354.2	.536 .551	0.0 334.7
i	.129	102.5	5-744	341.6	[]		1	.092	73.9	5.614	352.4 346.0		
i	.158 .192	125.8 152.7	5.629 5.062	339.5 336.3			ļ l	.127	101.4	5.562	345.7 342.1	-637	310.0
Ĭ	-25k	204.4	3.861	0.0			-	.153 .187	122.7	5.202 4.725	342.1	.659	287.8
- 1	285	229.9	4.503	356.8				.257 .289	204.2	3.129	335.9	1.487	192.0
- 1	.321	258.9	4.054	356.1				.289	229.3	3.884	358.0	1.044	176.3
-745	.032	26.7 52.1	7.230 6.933	349.8 346.6	705 664	336.0 332.7	-745			1		1	352.2
- 1	-086	74.1	6.693	341.4		332.1	رحان	.030	26.3	7.318	351 . 3 349 . 0	.744	337.1
1	ھي د.	102.0	6.192	341.1	-743	310.8	\	.058	50.5 75.8	6.674	343.7	-	
	.236	124.7 206.0	5.863 4.294	332.1	.613	312.6	1	.116	99.8	6.328	342.7	1.006	314.4
	.262	228.4	4.719	354.9 345.5				.146	126.0	5.942 4.875	337 · 3 355 · 4	.832	214.5
]	.291	253.3	4.248	343.6	.906	291.7		.273 .296	235.6 255.1	5.120 4.560	337.9 345.7	.838	210.6
.796	.029 .056	26.9 52.1	8.056 7.454	352.9 345.6	.881 .599	345.6 317.2	.798		1	1	354.4	ļ	1.5
1	.081	75.5	7.497	337.9			.190	.028	26.1 50.6	8.299 7.410	340.8	.777	322.0
1	•110	102.2	6.566 6.077	337.9 334.3	.765	303.4	Ι.	.079	73.6	7.030 6.380	339.8		
- 1	•137 •193	127.3	3.642	330.1 349.5				.109	101.3	6.380 5.297	331.4	.804	294.2
1	-221	207.3	4.442	SPR U	.6l/h	317.3		.135	176.5	4-123	335.8		
- 1	.250 .280		4.766	343.6 347.5	7	289.5	Ì	.221	205.3	4.573	345.3	.519	208.7
		262.5	5-340	- 1	1.177			279	233.9	5.128	325.7 339.8	.830	228.3
.825	.029	28.2	7.268	348.2 348.7	1.113	331.0 325.6	.827	~~	25.9	8.460	346.8		
- !	.078	75.6	6.764	340.2				.027	51.5	7.Qk5	344.1		
1	·100	102.7	6.302	331.3	-997	306.1		.079	76.1	7.528	1 328.0		
l	.130 .183	125.8 178.5	5.291 4.023	330.6 356.1	===			.107	103.4	7.528 6.133 2.884	330.2 345.7		: : :
- 1	.213	207.3	4.605	5.5	1.389	329.1		.210	203.3	3.100	2.7		- <i></i>
- 1	.238 .264	231.8 257.5	4.581 5.601	340.9 341.4	1.490	289.5		.240 .255	231.8	4.054	347.1 348.9] = = =
-857	.027	27.3	6.705	351.3	1.637	349.2	.860	.026	26.3	8.294	351.4	.997	355.2
-05(.051	51.1	6.214	345.8	1.757	338.9		051	51.6	7.660	334.9	.761	320.6
	.074	75.1	6.048	336.9				.074	74.3	6.404	324.6		283.4
l	.103	103.9	5.451 3.359	325.8 342.3	1.342	320.1 327.6		.098 ,153	99.2 151.2	3.707 3.976	315.4	.776 .643	298.6
- 1	.179	153.2 181.6	2,715	338.3				178	179.5	4.292	326.1		l
- 1	.202	204.9	3.483	331.7	1.518	352.4		-204	205.8	4.064	335.9	.770	183.2
- 1	229	232.2 256.5	3.292	333.6	2.161	307.4		.228 .252	230.4	4.321 3.773	317.3	1,606	148.8
.883	.027	27.8	6.494	350.5	1.702	0.6	.892	.025	26.1	9.005	340.4	-446	0.0
	.051	53.1	6.160	338.1	1.899	328.1		.049	50.3	7.821	333.4	1.172	209.3
- 1	.073	76.0	5.577 4.493	332.4 325.7	1.447	322.8		.073	75.3	6.752 3.753	317.0	.520	327.7
Ī	.146	152.1	3-493	336.2	1.247	335.1		-173	181.6	3.123	329.0		7.7.
1	.171	179.0 232.7	3.023	336.9 338.5				.202	212.2	3.685	332.2	.172	171.7
ı	248	259.6	3.928	338.7	1.951	309.5		-245	257.5	2.675	314.6	-755	322.4
مدو.	.116	125.7	3.369	343.7						1	1	Į	
- 1	.135	145.8 179.5	3.920	343.3	1.048	162.6		1	1	1	ĺ		1
	.197	213.0	4.358 4.484	319.9	1.004	174.9		1	1	1	1	1	1
- 1	.213	230.6	4.034	312.8	1			1	1	1	}	1	1
- 1	.240	259.5	3.133	299.7	1.007	173.0	<u> </u>	ــــــــــــــــــــــــــــــــــــــ	1	!	1		



MODEL PRESSURE-CELL LOCATIONS
[In Percent of Model Chord]

Cell number upper and lower surface	65A012 and 65A008	65A004 2-008, and 877A008
1 2 3 4 5 6 7 8 9 0 11 2 3 4 5 15	1.25 3.75 7.5 15 22.5 27.5 35 45 57.5 67.5 67.5 67.5 85 95	1.25 3.75 7.5 15 22.5 27.5 35 57.5 57.5 67.5 67.5 85 90



Figure 1.- Section profiles and pressure-cell locations of models.



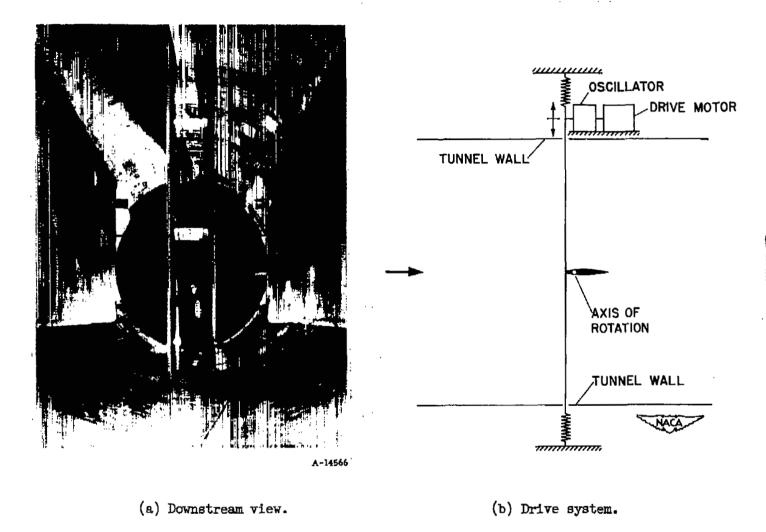


Figure 2.- View of test section with model in place and diagrammatic sketch of drive system.

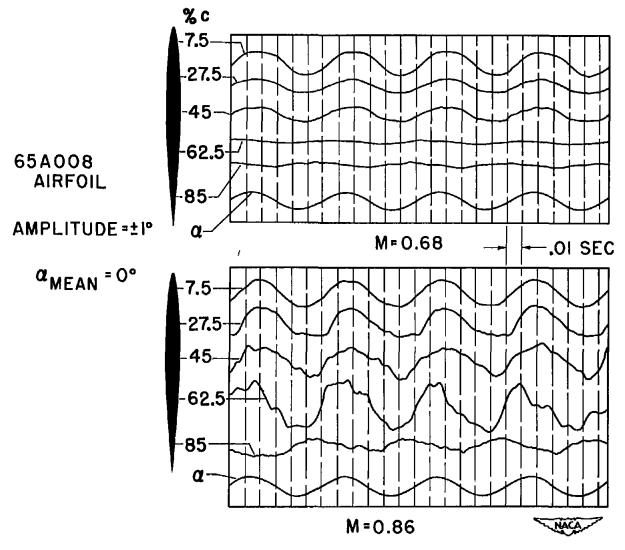


Figure 3.- Typical oscillograph traces.

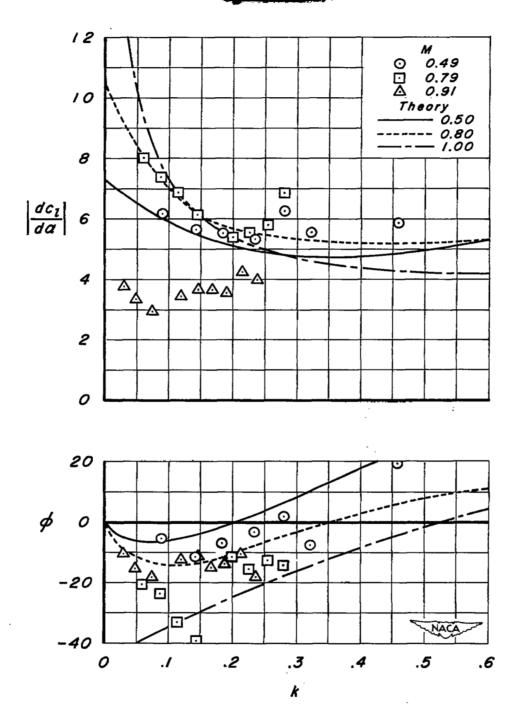
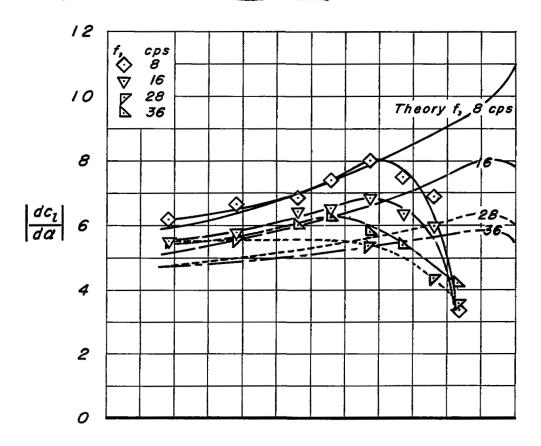


Figure 4.— Results as a function of reduced frequency, k, for several Mach numbers for the reference model, NACA 65A008; $\alpha_m = 0^{\circ}$.



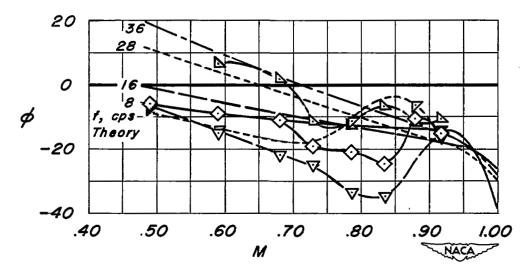
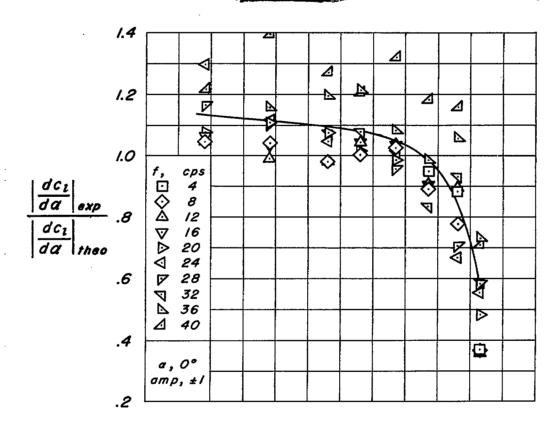


Figure 5.- Typical results for reference model, NACA 65A008; $\alpha_m = 0^\circ$.





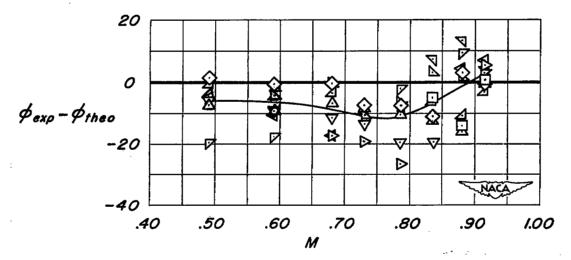
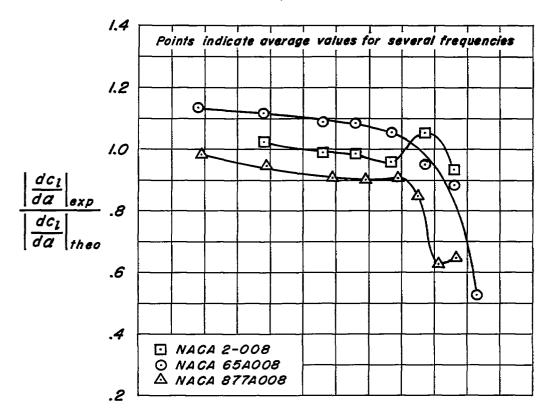
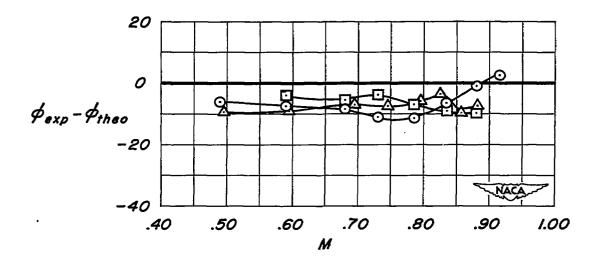


Figure 6 – Variation of experimental results from theory for reference model, NACA 65A008, with a faired line to show the mean variation with Mach number; $\alpha_m = 0^\circ$.



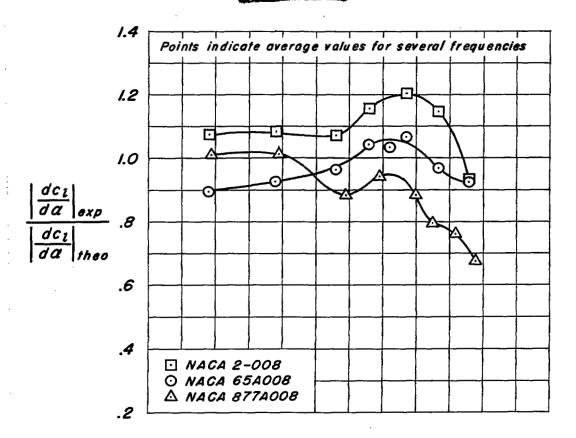


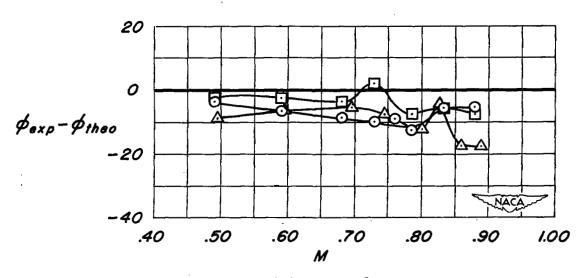




(a) am = 0°
Figure 7.- Effect of airfoil thickness distribution on lift derivatives.

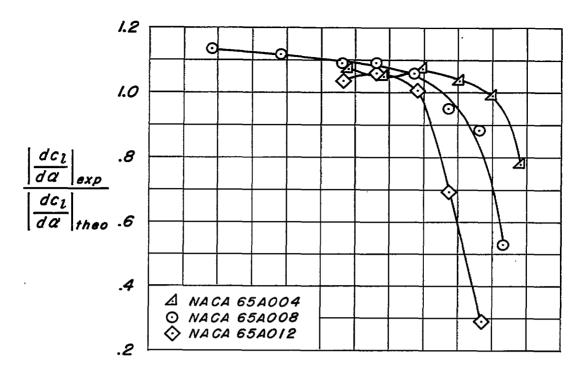






(b) $a_m = 2^{\circ}$ Figure 7.- Concluded.





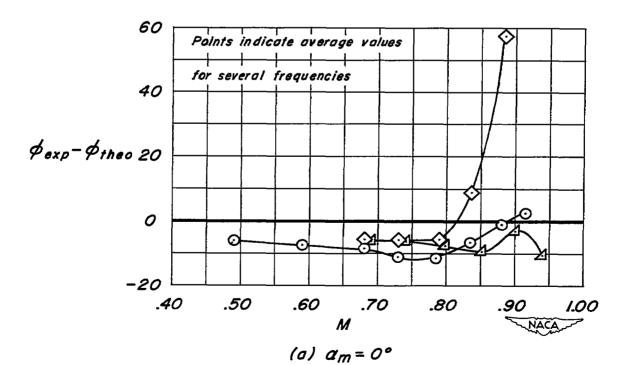
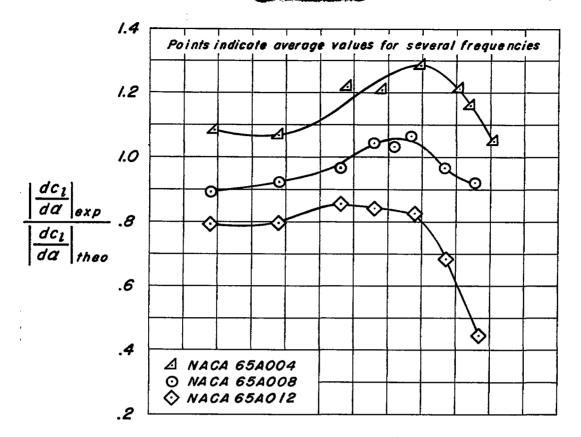
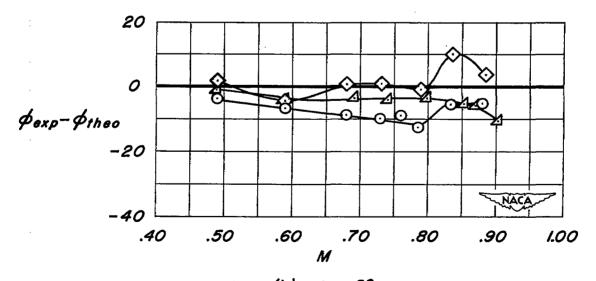


Figure 8.- Effect of airfoil thickness on lift derivatives.







(b) $\alpha_m = 2^{\circ}$ Figure 8.- Concluded.

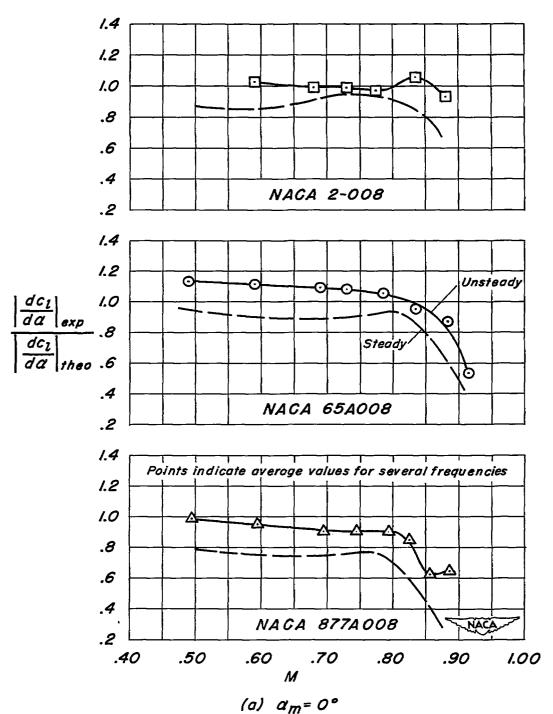
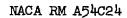
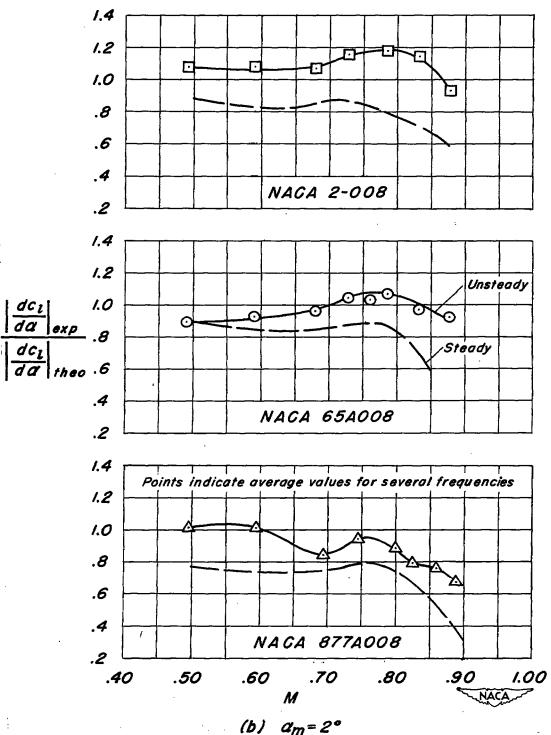
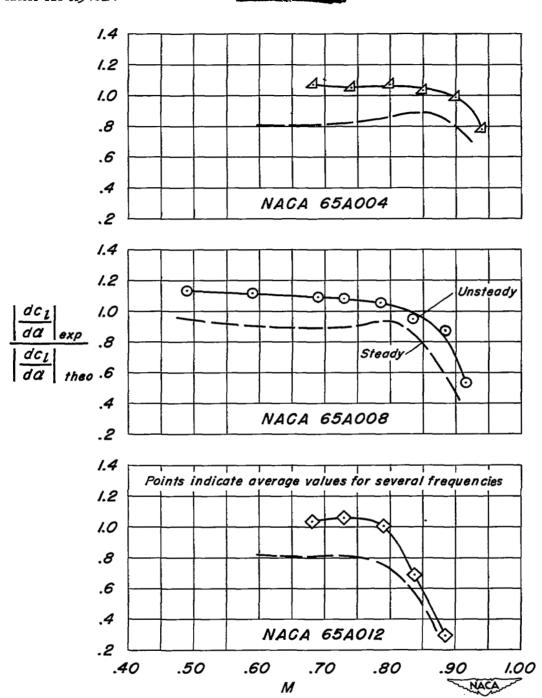


Figure 9.- Comparison of steady and unsteady lift derivatives for airfoils with varying thickness distributions.

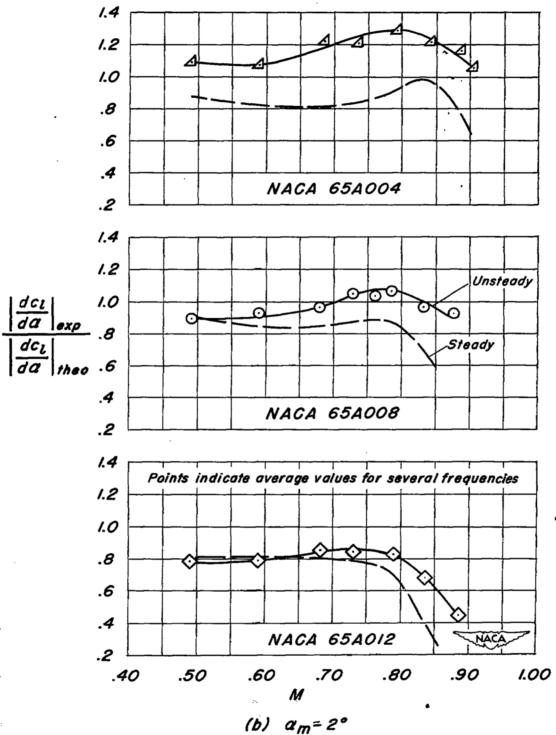




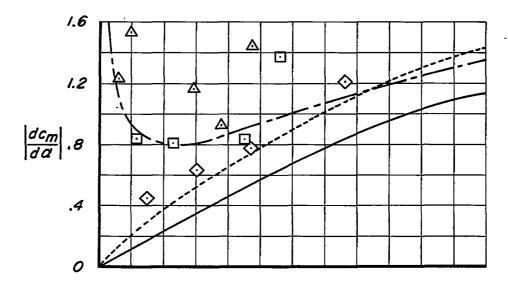
(D) $\alpha_m = 2^{\circ}$ Figure 9: Concluded.



(a) $\alpha_m = 0^\circ$ Figure 10: Comparison of steady and unsteady lift derivatives
for airfoils with varying thickness.



(b) $\alpha_m = 2^\circ$ Figure 10. – Concluded.



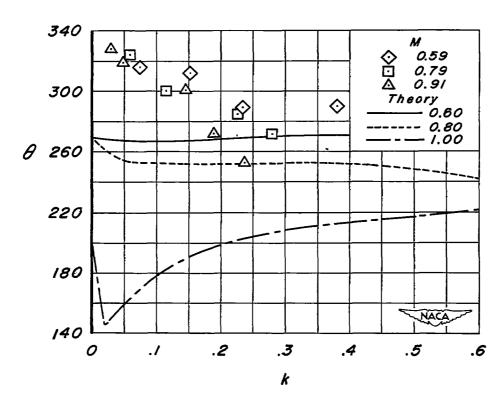


Figure II.- Results as a function of reduced frequency, k, for several Mach numbers for the reference model, NACA 65A008; $\alpha_m = 0^\circ$.



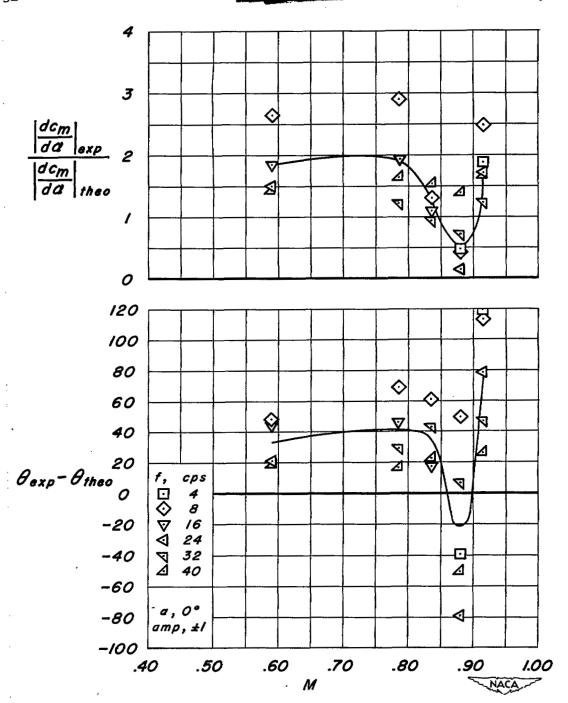
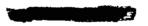
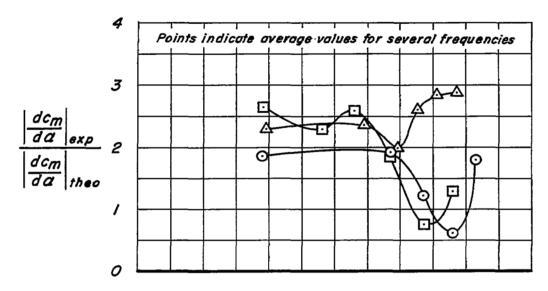


Figure 12.— Variation of experimental results from theory for reference model, NACA 65A008, with a faired line to show the mean variation with Mach number; $a_m = 0^\circ$.





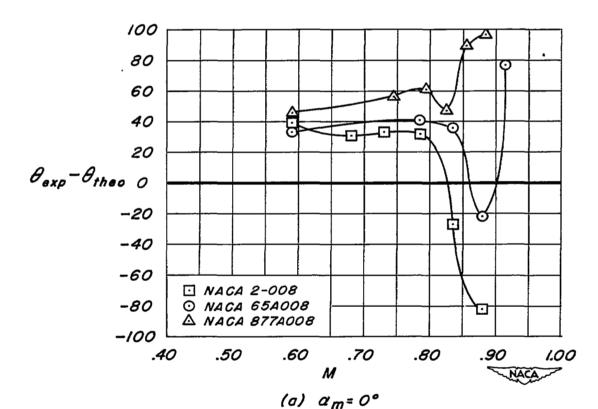
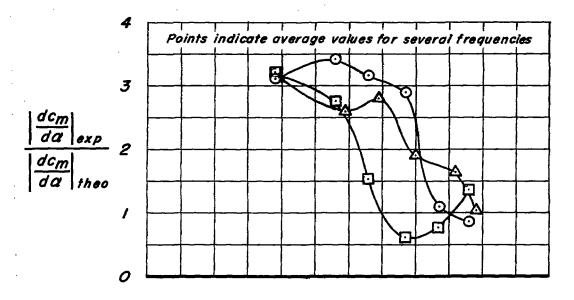
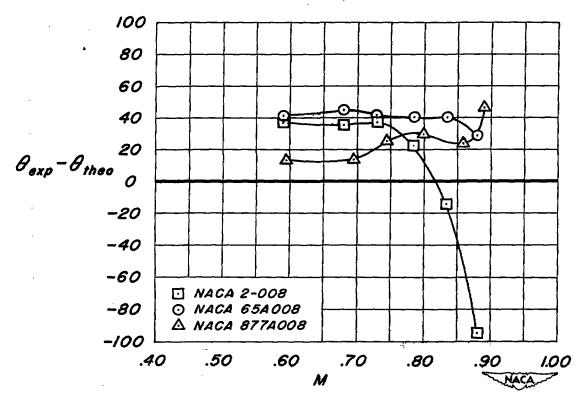


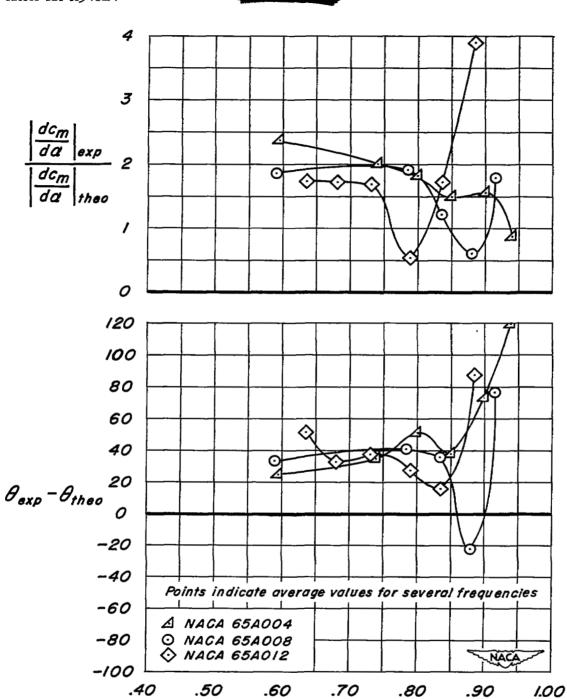
Figure 13.- Effect of airfoil thickness distribution on moment derivatives.





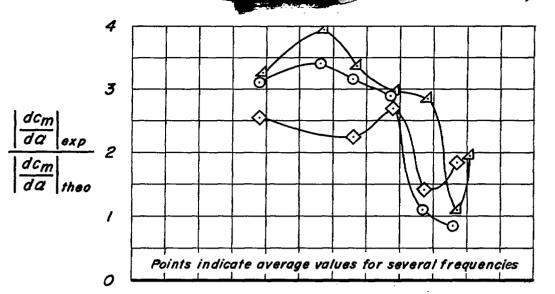


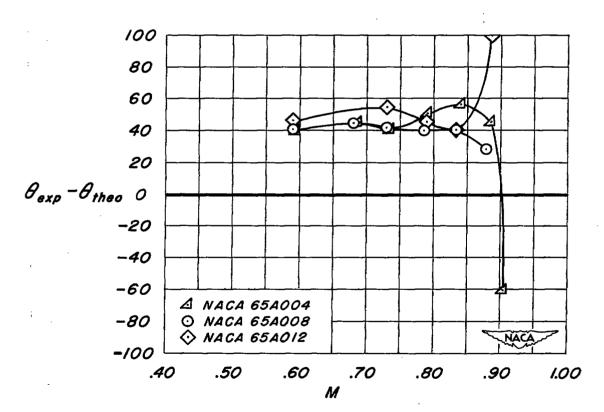
(b) $a_m = 2^{\circ}$ Figure 13.- Concluded.



(a) $\alpha_m = 0^\circ$ Figure 14.— Effect of airfoil thickness on moment derivatives.

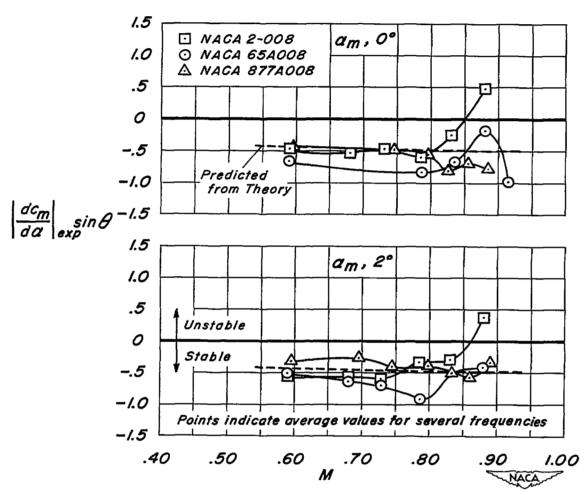




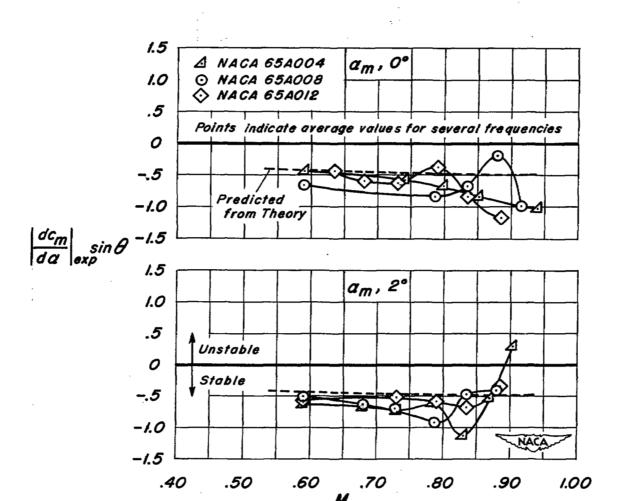


(b) $\alpha_m = 2^\circ$ Figure 14.- Concluded.





(a) Effect of airful thickness distribution.
Figure 15.— Damping component of the moment derivatives.



(b) Effect of airfoil thickness, Figure 15.- Concluded.